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Q^2 dependence of nuclear shadowing

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We reexamine the predictions of a two-phase model of shadowing in nuclear deep-inelastic scattering in light of new NMC data on the x , A , and Q^2 dependence of the ratios of structure functions. The model, which combines vector meson dominance at low Q^2 with diffractive Pomeron exchange for large Q^2 , agrees with the observed small, but nonzero, slopes in $\log Q^2$, which indicate the importance of higher twist shadowing effects in the transition region, $0.1 \leq Q^2 \leq 1 \text{ GeV}^2$. We note also that the latest E665 data on the deuteron to proton ratio suggests the presence of a small amount of shadowing in the deuteron.

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There has recently been a great deal of progress in the experimental study of the A dependence of nuclear structure functions in the small- x region where nuclear shadowing occurs. At CERN the New Muon Collaboration (NMC) has been able to extract the slopes in the four-momentum transfer squared, Q^2 , of the ratios of structure functions of a variety of nuclei for $x = Q^2/2M\nu$ down to ~ 0.001 [1–3], where ν is the energy transfer to the target and M the nucleon mass. Meanwhile, E665 at Fermilab has extended previous measurements on the deuteron and several heavier nuclei to x as low as $x \sim 10^{-5}$ [4–7].

These new measurements provide a tremendous challenge to theoretical models of nuclear deep-inelastic scattering. Not only do they cover five orders of magnitude in x , but the momentum transfer squared also varies from tens of GeV^2 to values as low as 0.05 GeV^2 . For the latter one is (to say the least) hard pressed to justify a parton model description. Clearly here one is dealing with a transition region between those regimes where hadronic and explicit quark treatments are appropriate. While it is unlikely that any single theoretical approach will be totally reliable over the entire range of x and Q^2 , it is important that all existing models be put to the test with a view to understanding the nature of the transition between high and low Q^2 .

Our aim is to extend earlier calculations of shadowing in deuterium [8] and heavier nuclei [9] to compare with these new data. The approach taken is a two-phase model, similar to the work of Kwiecinski and Badelek [10–12]. At high virtuality the interaction of the virtual photon with the nucleus is most efficiently parametrized in terms of diffractive scattering through the double and triple Pomeron. On the other hand, at low virtuality it is most natural to apply a vector meson dominance (VMD) model, in which the virtual photon interacts with the nucleons via its hadronic structure, namely the ρ^0 , ω , and ϕ mesons. The latter contribution vanishes at sufficiently high Q^2 , but in the region of interest here it is in fact responsible for the majority of the Q^2 variation.

For convenience we shall summarize the key ingredients of the calculation, beginning with the diffractive component. For this we use Pomeron (P) exchange between the projectile and two or more constituent nucleons to model the interaction of partons from different nucleons within the nucleus. Assuming factorization of the diffractive cross section [13], the shadowing correction (per nucleon) to the nuclear structure function F_2^A from P exchange is written as a convolution of the Pomeron structure function, F_2^P , with a distribution function (“flux factor”), $f_{P/A}$, describing the number density of exchanged Pomerons:

$$\delta^{(P)} F_2^A(x, Q^2) = \frac{1}{A} \int_{y_{\min}}^A dy f_{P/A}(y) F_2^P(x_P, Q^2), \quad (1)$$

where $y = x(1 + M_X^2/Q^2)$ is the light-cone momentum fraction carried by the Pomeron (M_X is the mass of the diffractive hadronic debris), and $x_P = x/y$ is the momentum fraction of the Pomeron carried by the struck quark in the Pomeron. The y dependence of $f_{P/A}(y)$ can be calculated within Regge theory [8–12, 14] (for a survey of other definitions of the “flux factor” see Ref. [15]). Within experimental errors, the factorization hypothesis, as well as the y dependence of $f_{P/A}(y)$ [8–12], are consistent with the results obtained recently by the H1 [16] and ZEUS [17] Collaborations at HERA from observations of large rapidity gap events in diffractive ep scattering. These data also confirm previous findings by the UA8 Collaboration [18] that the Pomeron structure function contains both a hard and a soft component: $F_2^P(x_P, Q^2) = F_2^{P(\text{hard})}(x_P, Q^2) + F_2^{P(\text{soft})}(x_P, Q^2)$. The hard component of F_2^P is generated from an explicit $q\bar{q}$ component of the Pomeron, and has an x_P dependence given by $x_P(1-x_P)$ [14, 19], in agreement with the recent diffractive data [16–18]. The soft part, which is driven at small x_P by the triple-Pomeron interaction [10, 19], has a sea quark-like x_P dependence,¹ with normalization fixed by the triple-Pomeron coupling constant [20].

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¹In Ref. [19] the soft component is associated with the $q\bar{q}g$ Fock component of the γ^* wave function.

The dependence of F_2^P on Q^2 at large Q^2 arises from radiative corrections to the parton distributions in the Pomeron [11,21], which leads to a weak, logarithmic, Q^2 dependence for the shadowing correction $\delta^{(P)}F_2^A$. Alone, the Pomeron contribution to shadowing would give a structure function ratio F_2^A/F_2^D that would be almost flat for $Q^2 \geq 2$ GeV². The low- Q^2 extrapolation of the $q\bar{q}$ component is parametrized by applying a factor $Q^2/(Q^2+Q_0^2)$, where $Q_0^2 \approx 0.485$ GeV² [22] may be interpreted as the inverse size of partons inside the virtual photon. For the nucleon sea quark densities relevant for $F_2^{P(\text{soft})}$ we use the recent parametrization from Ref. [22], which includes a low- Q^2 limit consistent with the real photon data, in which case the total Pomeron contribution $\delta^{(P)}F_2^A \rightarrow 0$ as $Q^2 \rightarrow 0$.

To adequately describe shadowing for small Q^2 requires one to use a higher-twist mechanism, such as vector meson dominance. VMD is empirically based on the observation that some aspects of the interaction of photons with hadronic systems resemble purely hadronic interactions [23,24]. In terms of QCD this is understood in terms of a coupling of the photon to a correlated $q\bar{q}$ pair of low invariant mass, which may be approximated as a virtual vector meson. One can

then estimate the amount of shadowing in terms of the multiple scattering of the vector meson using Glauber theory. The corresponding correction (per nucleon) to the nuclear structure function is

$$\delta^{(V)}F_2^A(x, Q^2) = \frac{1}{A} \frac{Q^2}{\pi} \sum_V \frac{M_V^4 \delta\sigma_{VA}}{f_V^2(Q^2 + M_V^2)^2}, \quad (2)$$

where $\delta\sigma_{VA}$ is the shadowing correction to the vector meson–nucleus cross section, calculated in Ref. [9], f_V is the photon–vector meson coupling strength [23] (see also [25]), and M_V is the vector meson mass. In practice, only the lowest mass vector mesons ($V = \rho^0, \omega, \phi$) are important at low Q^2 . (Inclusion of higher mass vector mesons, including continuum contributions, leads to so-called generalized vector meson dominance models [26].) Usually one omits nondiagonal vector meson transitions ($VN \rightarrow V'N$), as these are not expected to be large. The vector meson propagators in Eq. (2) lead to a strong Q^2 dependence of $\delta^{(V)}F_2^A$, which peaks at $Q^2 \sim 1$ GeV². For $Q^2 \rightarrow 0$ and fixed x , $\delta^{(V)}F_2^A$ disappears due to the vanishing of the total F_2^A . Furthermore, since this is a higher twist effect, shadowing in the VMD

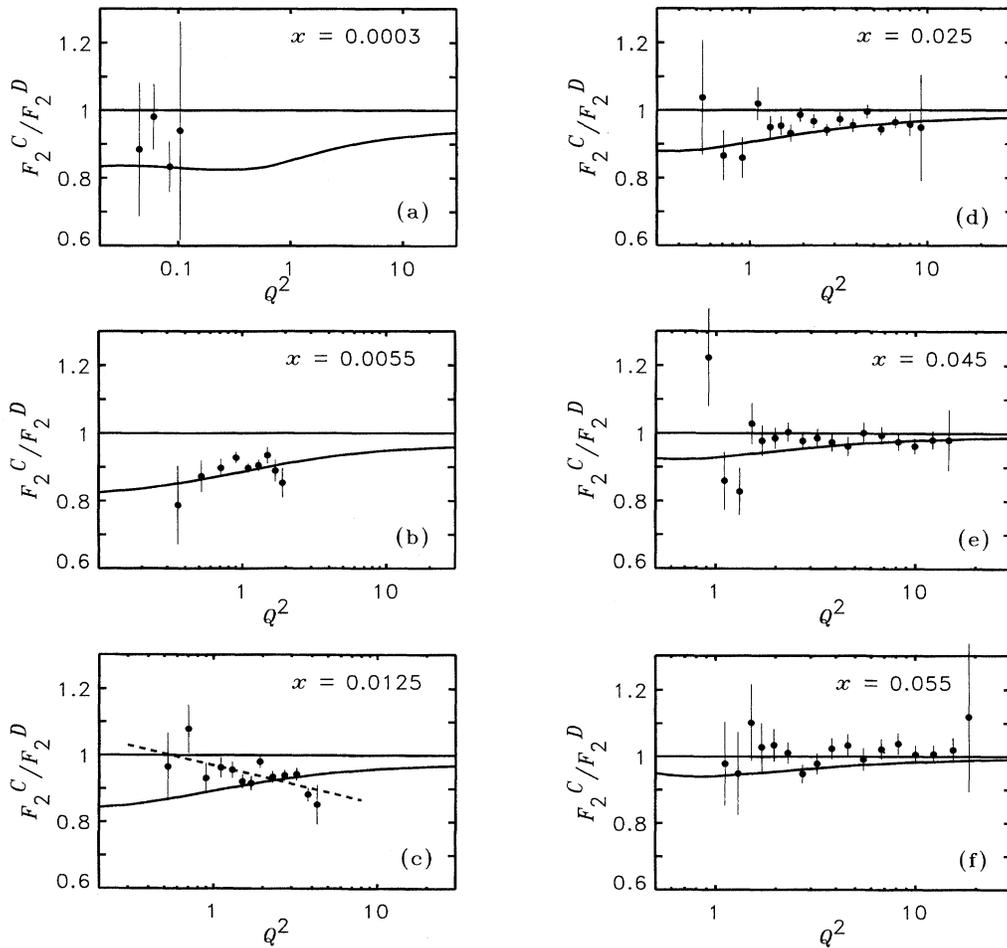


FIG. 1. Q^2 variation of the C/D structure function ratio, for various values of x ($x = 0.0003, 0.0055, 0.0125, 0.025, 0.045,$ and 0.055), compared with the complete NMC data [2,3]. Also shown for $x=0.0125$ is a linear fit in $\log Q^2$ to the data.

model dies off quite rapidly between $Q^2 \sim 1$ and 10 GeV^2 , so that for $Q^2 \geq 10 \text{ GeV}^2$ it is almost negligible—leaving only the diffractive term, $\delta^{(p)} F_2^A$. (Note that at fixed ν , for decreasing Q^2 the ratio F_2^A/F_2^D approaches the photoproduction limit.)

While the asymptotic Q^2 behavior of nuclear shadowing seems clear, there is still considerable interest in the transition region where the high- and low- Q^2 descriptions merge. In practice, this occurs for Q^2 between about 0.5 and 5 GeV^2 , which is precisely the region where most of the recent NMC data on C, Ca, and other nuclei have been taken [1–3].

For light nuclei, such as C, the dominant mechanism for nuclear shadowing involves the double scattering of the projectile from two nucleons. Higher order terms (multiple re-scattering) in the Glauber expansion attenuate the incident flux of vector mesons (or of the hadronic state X for the diffractive component [11]) as they traverse the nucleus, which will be progressively more important as A increases [9]. For the VMD component, the magnitude of the attenuation is determined by the mean free path, L_V , of the vector meson in the nucleus, $L_V = (\rho \sigma_{VN})^{-1}$, where σ_{VN} is the total VN cross section [27]. If the mean free path $L_X = (\rho \sigma_{XN})^{-1}$ of the hadronic state X is independent of the

mass M_X [11], one may take the reabsorption cross section $\sigma_{XN} \sim \sigma_{VN} \sim 20\text{--}30 \text{ mb}$, although in the Q^2 range covered by the NMC data the F_2^A/F_2^D ratios are not very sensitive to the precise value of σ_{XN} . For the single particle density ρ in heavy nuclei ($A \geq 16$) we use the Woods-Saxon (or Fermi) density, while for light nuclei ($A \leq 16$) the harmonic oscillator (shell model) form is more appropriate [28]. Short range correlations are included through a Fermi gas correlation function which puts a “hole” in the two-body density approximately 0.5 fm wide at $1/2$ maximum density [9]. The inclusion of correlations has the effect of decreasing slightly the amount of shadowing (i.e., increasing the ratio F_2^A/F_2^D) at low Q^2 .

Having outlined the essential features of the model, we now turn to a detailed comparison with the data. In Fig. 1 we plot the ratio of the structure functions F_2 (normalized to one nucleon) for C and D as a function of Q^2 , for various values of x ranging from $x = 0.0003$ to $x = 0.055$. The data represent the complete sample taken by the NMC for C nuclei. The overall agreement between the model calculation and the data is clearly excellent. In particular, the observed Q^2 dependence of the ratios is certainly compatible with that indicated by the NMC data. At large Q^2 ($Q^2 \geq 10 \text{ GeV}^2$), the

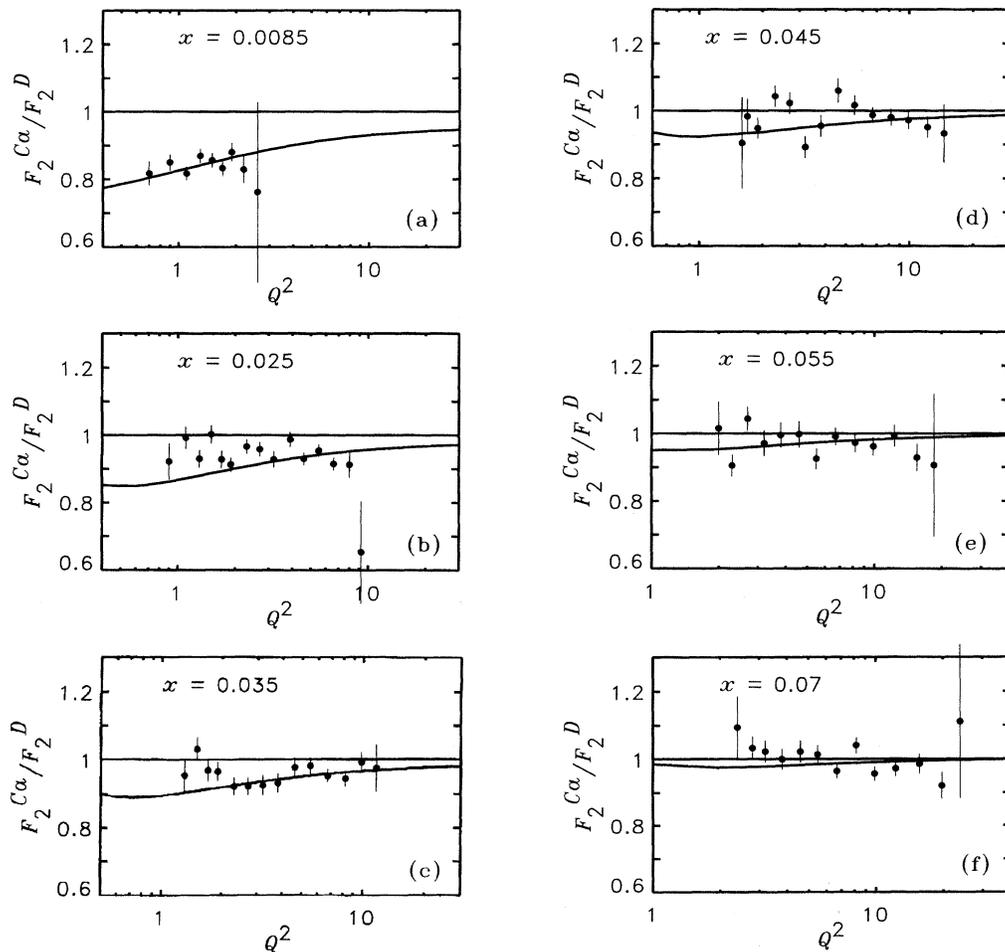


FIG. 2. Q^2 variation of the Ca/D structure function ratio, for $x = 0.0085, 0.025, 0.035, 0.045, 0.055,$ and 0.07 , compared with the NMC data [2].

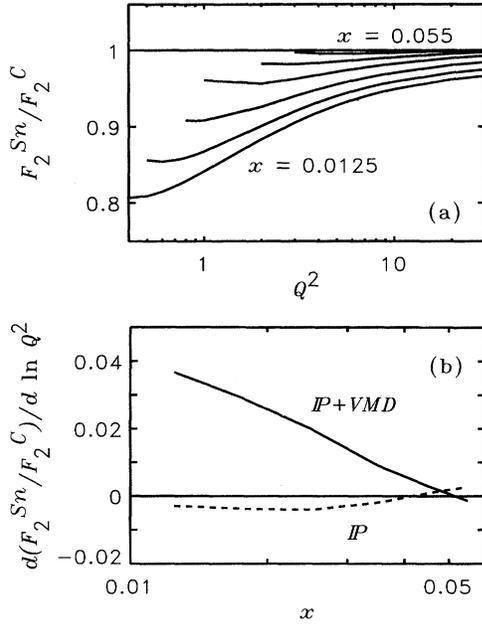


FIG. 3. (a) Model prediction for the Q^2 dependence of the structure function ratio of Sn to C, for $x = 0.0125$ (lowest curve), 0.0175, 0.025, 0.035, 0.045, and 0.055 (highest curve). (b) Slope in $\log Q^2$ of the Sn/C ratio as a function of x —solid curve is the full result, dashed is the Pomeron contribution only.

calculated curves are almost constant with Q^2 , as would be expected from a partonic (scaling) mechanism [19,29–32]. However, in the smallest x bins the Q^2 values reach as low as $Q^2 \approx 0.05 \text{ GeV}^2$. Clearly this region of Q^2 is inaccessible to any model involving only a partonic mechanism, and it is essential here to invoke a nonscaling mechanism such as the VMD model.

It is common practice in many data analyses to extract slopes in $\log Q^2$ by performing simple straight-line fits to the data—such as $F_2^A/F_2^D = a + b \log Q^2$. To illustrate a potential difficulty with using such extracted slopes to discriminate between different models of shadowing, we plot in Fig. 1 (for $x = 0.0125$) the slope that one would obtain by making a linear fit to the $\log Q^2$ variation of F_2^C/F_2^D . Evidently the slope thus obtained is negative, while the calculated curves have a positive slope. Within the present degree of accuracy, the C/D data are equally compatible with a positive or negative slope, and the trend of the data as a function of Q^2 clearly favors the positive value obtained in our calculation.

In Fig. 2 we show the structure function ratio for Ca/D, for x between $x = 0.0085$ and 0.07. Again, within the error bars, the calculated x and Q^2 dependence is in excellent agreement with the data. The apparent small decrease in the ratio for the largest Q^2 points could be due to a small nuclear dependence of the ratio of the longitudinal to transverse virtual photoabsorption cross sections [33], which is assumed to be zero in the data extraction. However, the statistics on the data do not allow a definitive statement about this effect.

In the future the NMC will also produce data on the Q^2 dependence of the Sn to C ratio [34]. In Fig. 3(a) we show the predictions for the structure function ratio as a function of Q^2 , for $x = 0.0125$ (lowest curve), 0.0175, 0.025, 0.035,

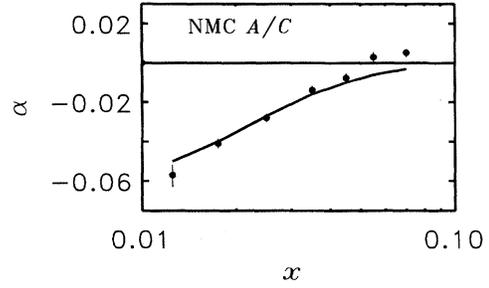


FIG. 4. x dependence of the slope α from the structure function ratio $F_2^A/F_2^C \propto A^\alpha$, compared with NMC data on $A = D, \text{Li, Be, Al, Ca, Fe, and Sn}$.

0.045 and 0.055 (highest curve). Because of the expected higher accuracy of these data [34], linear fits to the $\log Q^2$ dependence should result in more reliable extracted slopes. In Fig. 3(b) we show the slopes $b = d(F_2^{Sn}/F_2^C)/d \log Q^2$, extracted from the curves in Fig. 3(a), as a function of x . The solid line is the slope for the VMD and P-exchange mechanisms, while the dashed line represents the Pomeron contribution only. We would like to stress that essentially zero slopes would be obtained in models where only a partonic mechanism would be utilized [19,29–32]. Observation of nonzero slopes would clearly support the hypothesis that the intermediate- Q^2 region is dominated by nonscaling, higher-twist effects, having a nontrivial Q^2 dependence.

Data for other nuclei have also been taken recently by the NMC. The A dependence of the structure function ratio of deuterium, Li, Be, Al, Ca, Fe, and Sn to carbon has been parametrized as $F_2^A(x)/F_2^C(x) \propto A^{\alpha(x)}$. In Fig. 4 we show the slope $\alpha(x) = d(F_2^A/F_2^C)/d \log A$ as a function of x . The agreement at the low x values is clearly very good. At larger values of x ($x \geq 0.07$) the data go above unity, and tend to lie slightly above the calculated curve. This is a reflection of the fact that the present model has no mechanism for antishadowing built in.

Having obtained excellent fits to the nuclear data, we can be reasonably confident that applying the same model to the deuteron [8,12,35,36] will yield reliable results. Indeed, here one has fewer model parameters to deal with, since all of the shadowing is generated through double scattering alone. As was observed in Refs. [8,12], the largest uncertainty in the calculation of the shadowing in deuterium is the deuteron

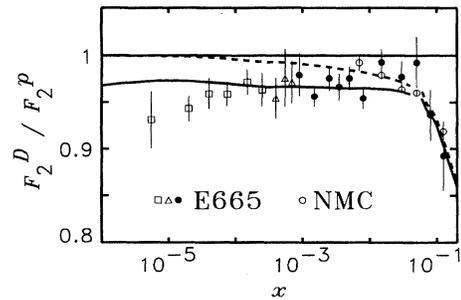


FIG. 5. x dependence of the D/p structure function ratio, compared with the low- x E665 data [4] and NMC data [37] at larger x . The dashed curve is the result without any shadowing correction.

wave function. Nevertheless, the presence of shadowing in the deuteron would be confirmed through observation of a deviation from unity in the D/p structure function ratio in the kinematic region where Regge theory is expected to be valid. Although the exact value of x below which the proton and (free) neutron structure functions become equivalent is not known, it is expected that at low enough x , $F_2^p \rightarrow F_2^n$, in which case $F_2^D/F_2^p \rightarrow 1 + \delta F_2^D/F_2^p$. In Fig. 5 we show the data at very low x taken by the E665 Collaboration [4], as well as the earlier NMC data at larger x [37]. The calculated ratio with a small shadowing correction is shown by the solid curve, while the result for the case of no shadowing is indicated by the dashed curve. The data clearly favor the shadowing scenario. As a fraction of F_2^D , the shadowing correction amounts to about 1.5% at $x = 10^{-2}$ up to about 3% for $x \leq 10^{-5}$.

In summary, we have seen that the latest data from the

NMC and E665 indicate that nuclear shadowing in the low- and intermediate- Q^2 regions is controlled by the dynamics of a higher-twist, vector meson dominance mechanism. At larger Q^2 this component disappears, leaving behind a scaling component which is understood to arise from diffractive scattering from the Pomeron component of the nucleon, and which agrees with the approximate Q^2 independence of the data at large Q^2 . Models based solely on partonic mechanisms can therefore provide only limited insight into the physics of nuclear shadowing.

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- [1] NM Collaboration, P. Amaudruz *et al.*, *Z. Phys. C* **51**, 387 (1991).
- [2] NM Collaboration, P. Amaudruz *et al.*, *Nucl. Phys.* **B441**, 3 (1995).
- [3] NM Collaboration, M. Arneodo *et al.*, *Nuovo Cimento* **107A**, 2141 (1994).
- [4] E665 Collaboration, M.R. Adams *et al.*, *Phys. Rev. Lett.* **75**, 1466 (1995).
- [5] E665 Collaboration, M.R. Adams *et al.*, *Phys. Lett. B* **309**, 477 (1993).
- [6] E665 Collaboration, M.R. Adams *et al.*, Fermilab Report No. hep-ex/9505006, April 1995.
- [7] E665 Collaboration, M.R. Adams *et al.*, *Phys. Rev. Lett.* **68**, 3266 (1992).
- [8] W. Melnitchouk and A.W. Thomas, *Phys. Rev. D* **47**, 3783 (1993).
- [9] W. Melnitchouk and A.W. Thomas, *Phys. Lett. B* **317**, 437 (1993).
- [10] J. Kwiecinski and B. Badelek, *Phys. Lett. B* **208**, 508 (1988).
- [11] J. Kwiecinski, *Z. Phys. C* **45**, 461 (1990).
- [12] B. Badelek and J. Kwiecinski, *Nucl. Phys.* **B370**, 278 (1992); *Phys. Rev. D* **50**, 4 (1994).
- [13] G. Ingelman and P.E. Schlein, *Phys. Lett.* **152B**, 256 (1985).
- [14] A. Donnachie and P.V. Landshoff, *Phys. Lett. B* **191**, 309 (1987).
- [15] K. Prytz, *Z. Phys. C* **64**, 79 (1994).
- [16] H1 Collaboration, T. Ahmed *et al.*, *Phys. Lett. B* **348**, 681 (1995).
- [17] ZEUS Collaboration, M. Derrick *et al.*, DESY Report No. 95-093, 1995.
- [18] UA8 Collaboration, A. Brandt *et al.*, *Phys. Lett. B* **297**, 417 (1992).
- [19] N.N. Nikolaev and B.G. Zakharov, *Z. Phys. C* **49**, 607 (1991); **53**, 331 (1992).
- [20] K. Goulianos, *Phys. Rep.* **101**, 169 (1983); Rockefeller U. Report No. RU 95/E-06, 1995.
- [21] J.C. Collins, J. Huston, J. Pumplin, H. Weerts, and J.J. Whitmore, *Phys. Rev. D* **51**, 3182 (1995).
- [22] A. Donnachie and P.V. Landshoff, *Z. Phys. C* **61**, 139 (1994).
- [23] T.H. Bauer, R.D. Spital, D.R. Yennie, and F.M. Pipkin, *Rev. Mod. Phys.* **50**, 261 (1978).
- [24] G.A. Schuler and T. Sjöstrand, *Nucl. Phys.* **B407**, 539 (1993); *Phys. Rev. D* **49**, 2257 (1994).
- [25] H.B. O'Connell *et al.*, *Phys. Lett. B* **336**, 1 (1994); H.B. O'Connell *et al.*, in *Trends in Particle and Nuclear Physics*, edited by W.-Y. Pauchy Hwang (Plenum Press) (submitted).
- [26] C.L. Bilchak, D. Schildknecht, and J.D. Stroughair, *Phys. Lett. B* **214**, 441 (1988); **233**, 461 (1989); L.L. Frankfurt and M.I. Strikman, *Phys. Rep.* **160**, 235 (1988); *Nucl. Phys.* **B316**, 340 (1989); G. Piller and W. Weise, *Phys. Rev. C* **42**, R1834 (1991); G. Piller, W. Ratzka, and W. Weise, *Z. Phys. A* (to be published); G. Shaw, *Phys. Lett. B* **228**, 125 (1989); *Phys. Rev. D* **47**, R3676 (1993).
- [27] A. Donnachie and P.V. Landshoff, *Phys. Lett. B* **296**, 227 (1992).
- [28] G. Grammar and J.D. Sullivan, in *Electromagnetic Interactions of Hadrons (Vol. 2)*, edited by A. Donnachie and G. Shaw (Plenum Press, New York, 1978).
- [29] A.H. Mueller and J. Qiu, *Nucl. Phys.* **B268**, 427 (1986); J. Qiu, *ibid.* **B291**, 746 (1987); E.L. Berger and J. Qiu, *Phys. Lett. B* **206**, 141 (1988); F.E. Close, J. Qiu, and R.G. Roberts, *Phys. Rev. D* **40**, 2820 (1989).
- [30] S.J. Brodsky and H.J. Lu, *Phys. Rev. Lett.* **64**, 1342 (1990).
- [31] S. Kumano, *Phys. Rev. C* **48**, 2016 (1993); **50**, 1247 (1994); R. Kobayashi, S. Kumano, and M. Miyama, Saga Report No. SAGA-HE-78-95, 1995.
- [32] B. Kopeliovich and B. Povh, Heidelberg Report No. MPIH-V12-1995, 1995.
- [33] J. Milana, *Phys. Rev. C* **49**, 2820 (1994).
- [34] A. Mücklich (private communication).
- [35] V.R. Zoller, *Z. Phys. C* **54**, 425 (1992).
- [36] H. Khan and P. Hoodbhoy, *Phys. Lett. B* **298**, 181 (1993).
- [37] NM Collaboration, M. Arneodo *et al.*, *Phys. Rev. D* **50**, 1 (1994).